

# Dark Matter Annihilating into Sterile Neutrino

Yi-Lei Tang

Center for High Energy Physics, Peking University

October 26, 2016

- ▶ This slide is based on JHEP 1603 (2016) 043, Yi-Lei Tang, Shou-Hua Zhu, and arXiv:1609.07841.

# Introduction to Seesaw Mechanisms

- ▶ Right-handed neutrino  $N$  can result in the light neutrinos' masses through Type-I seesaw mechanisms.

# Introduction to Seesaw Mechanisms

- ▶ Right-handed neutrino  $N$  can result in the light neutrinos' masses through Type-I seesaw mechanisms.
- ▶ Majorana mass among right-handed neutrinos.

# Introduction to Seesaw Mechanisms

- ▶ Right-handed neutrino  $N$  can result in the light neutrinos' masses through Type-I seesaw mechanisms.
- ▶ Majorana mass among right-handed neutrinos.
- ▶

$$M = \begin{bmatrix} 0 & m_D \\ m_D^T & m_N \end{bmatrix} \quad (1)$$

# Introduction to Seesaw Mechanisms

- ▶ Right-handed neutrino  $N$  can result in the light neutrinos' masses through Type-I seesaw mechanisms.
- ▶ Majorana mass among right-handed neutrinos.
- ▶

$$M = \begin{bmatrix} 0 & m_D \\ m_D^T & m_N \end{bmatrix} \quad (1)$$

- ▶  $\rightarrow m_\nu = -\frac{m_D^2}{m_N}$ .

# Introduction to Seesaw Mechanisms

- ▶ Right-handed neutrino  $N$  can result in the light neutrinos' masses through Type-I seesaw mechanisms.
- ▶ Majorana mass among right-handed neutrinos.
- ▶

$$M = \begin{bmatrix} 0 & m_D \\ m_D^T & m_N \end{bmatrix} \quad (1)$$

- ▶  $\rightarrow m_\nu = -\frac{m_D^2}{m_N}$ .
- ▶  $m_D = y_\nu v_{\text{EW}}$ , usually  $y \sim 1$ , and  $m_N \gg 1$  TeV.

# Introduction to Seesaw Mechanisms

- ▶ Right-handed neutrino  $N$  can result in the light neutrinos' masses through Type-I seesaw mechanisms.
- ▶ Majorana mass among right-handed neutrinos.
- ▶

$$M = \begin{bmatrix} 0 & m_D \\ m_D^T & m_N \end{bmatrix} \quad (1)$$

- ▶  $\rightarrow m_\nu = -\frac{m_D^2}{m_N}$ .
- ▶  $m_D = y_\nu v_{EW}$ , usually  $y \sim 1$ , and  $m_N \gg 1$  TeV.
- ▶  $y_\nu \sim 10^{-7}-10^{-5}$ ,  $m_N < 1$  TeV (Naive TeV Seesaw).



# Introduction to Seesaw Mechanisms

- ▶ Right-handed neutrino  $N$  can result in the light neutrinos' masses through Type-I seesaw mechanisms.
- ▶ Majorana mass among right-handed neutrinos.
- ▶

$$M = \begin{bmatrix} 0 & m_D \\ m_D^T & m_N \end{bmatrix} \quad (1)$$

- ▶  $\rightarrow m_\nu = -\frac{m_D^2}{m_N}$ .
- ▶  $m_D = y_\nu v_{EW}$ , usually  $y \sim 1$ , and  $m_N \gg 1$  TeV.
- ▶  $y_\nu \sim 10^{-7}-10^{-5}$ ,  $m_N < 1$  TeV (Naive TeV Seesaw).
- ▶ For linear see-saw or inverse see-saw (pseudo-Dirac sterile neutrino),  $y_\nu$  can be as large as  $10^{-3}$ .

# Sterile Neutrino-Portal Dark Matter?

- ▶ Sterile Neutrino Dark Matter? ( $m_{\text{DM}} \ll 1 \text{ GeV}$ , “Warm Dark Matter”)

# Sterile Neutrino-Portal Dark Matter?

- ▶ Sterile Neutrino Dark Matter? ( $m_{\text{DM}} \ll 1 \text{ GeV}$ , “Warm Dark Matter”)
- ▶ What about the dark matter annihilate into sterile neutrinos?

# Sterile Neutrino-Portal Dark Matter?

- ▶ Sterile Neutrino Dark Matter? ( $m_{\text{DM}} \ll 1 \text{ GeV}$ , “Warm Dark Matter”)
- ▶ What about the dark matter annihilate into sterile neutrinos?
- ▶ Two previous examples,

# Sterile Neutrino-Portal Dark Matter?

- ▶ Sterile Neutrino Dark Matter? ( $m_{\text{DM}} \ll 1 \text{ GeV}$ , “Warm Dark Matter”)
- ▶ What about the dark matter annihilate into sterile neutrinos?
- ▶ Two previous examples,
- ▶ 1) NMSSM+Right-handed Neutrino.

# Sterile Neutrino-Portal Dark Matter?

- ▶ Sterile Neutrino Dark Matter? ( $m_{\text{DM}} \ll 1 \text{ GeV}$ , “Warm Dark Matter”)
- ▶ What about the dark matter annihilate into sterile neutrinos?
- ▶ Two previous examples,
- ▶ 1) NMSSM+Right-handed Neutrino.
- ▶  $\text{DM} + \text{DM} \rightarrow N + N$  usually does not dominate, but can be important sometimes.

# Sterile Neutrino-Portal Dark Matter?

- ▶ Sterile Neutrino Dark Matter? ( $m_{\text{DM}} \ll 1 \text{ GeV}$ , “Warm Dark Matter”)
- ▶ What about the dark matter annihilate into sterile neutrinos?
- ▶ Two previous examples,
  - ▶ 1) NMSSM+Right-handed Neutrino.
  - ▶  $\text{DM} + \text{DM} \rightarrow N + N$  usually does not dominate, but can be important sometimes.
  - ▶ 2) MSSM+(B-L) $Z'$ , R. Allahverdi, et.al., 0907.1486, etc.,

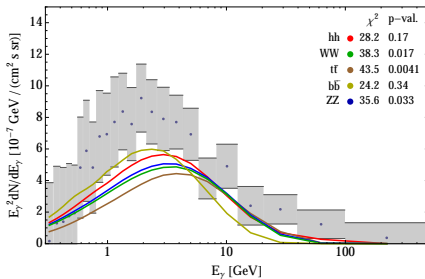
# Sterile Neutrino-Portal Dark Matter?

- ▶ Sterile Neutrino Dark Matter? ( $m_{\text{DM}} \ll 1 \text{ GeV}$ , “Warm Dark Matter”)
- ▶ What about the dark matter annihilate into sterile neutrinos?
- ▶ Two previous examples,
  - ▶ 1) NMSSM+Right-handed Neutrino.
    - ▶  $\text{DM} + \text{DM} \rightarrow N + N$  usually does not dominate, but can be important sometimes.
  - ▶ 2) MSSM+(B-L) $Z'$ , R. Allahverdi, et.al., 0907.1486, etc.,
    - ▶ In this model, there are some parameter space that  $\text{DM} + \text{DM} \rightarrow N + N$  can dominate.



# Current fittings on the galactic center excess of the $\gamma$ -ray

- ▶ DM  $\rightarrow b\bar{b}$  fits the **galactic center excess (GCE)** well.  
 $W^+W^-$ ,  $ZZ$ ,  $t\bar{t}$  do not.



**Figure** : From arXiv:1411.2592, by Prateek Agrawal, Brian Batell, Patrick J. Fox, and Roni Harnik. Data from F. Calore, et.al., 1409.0042.

# Current fittings on the galactic center excess of the $\gamma$ -ray

- DM  $\rightarrow b\bar{b}$  fits the galactic center excess (GCE) well.  
 $W^+W^-$ ,  $ZZ$ ,  $t\bar{t}$  do not.

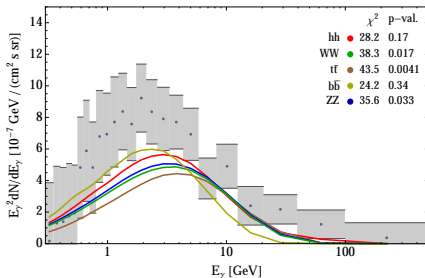


Figure : From arXiv:1411.2592, by Prateek Agrawal, Brian Batell, Patrick J. Fox, and Roni Harnik. Data from F. Calore, et.al., 1409.0042.

- The key is the position of the peak and the length of the tail!  
 $W/Z/t$  is too heavy for a lighter peak.

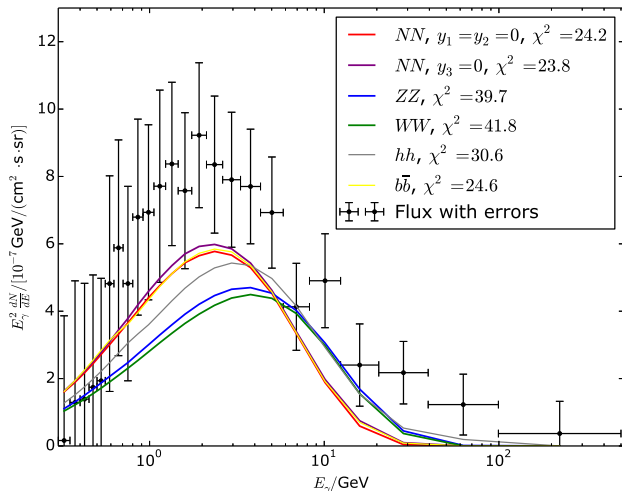
# Model Independent Analysis of the GCE Originating From $DM+DM \rightarrow N + N$

- ▶  $DM+DM \rightarrow N + N$ , RHN  $\rightarrow$  off-shell  $W, Z$ , which might move the position of the peak downward.

# Model Independent Analysis of the GCE Originating From $DM+DM \rightarrow N + N$

- ▶  $DM+DM \rightarrow N + N$ , RHN  $\rightarrow$  off-shell  $W, Z$ , which might move the position of the peak downward.
- ▶ The best-fitted points are  $m_N = 32.0$  GeV,  $m_\chi = 44.2$  GeV, with  $\chi^2 = 24.22$  and the best-fitted  $\langle\sigma v\rangle = 2.63 \times 10^{-26}$  cm<sup>3</sup>/s for the  $y_1 = y_2 = 0, y_3 \neq 0$  case, and  $m_N = 27.0$  GeV,  $m_\chi = 45.4$  GeV, with  $\chi^2 = 23.81$  and the best-fitted  $\langle\sigma v\rangle = 3.37 \times 10^{-26}$  cm<sup>3</sup>/s for the  $y_3 = 0, y_1^2 + y_2^2 \neq 0$  case.

# Model Independent Analysis of the GCE Originating From $DM+DM \rightarrow N + N$



# Model Independent Analysis of the GCE Originating From $DM+DM \rightarrow N + N$

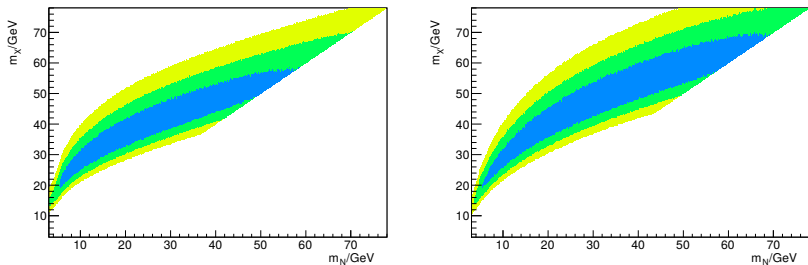


Figure : 1,2,3- $\sigma$  area fitting the GCE data

# Model Independent Analysis of the GCE Originating From $DM+DM \rightarrow N + N$

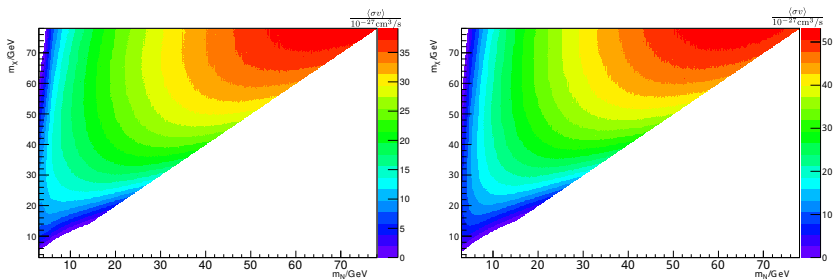


Figure : Best-fitted  $\langle \sigma v \rangle$  for the  $\gamma$ -ray GCE

# Standard WIMP Calculation?

- ▶ Standard WIMP calculation requires the annihilation products to fall into thermal equilibrium rapidly with the thermal bath.



# Standard WIMP Calculation?

- ▶ Standard WIMP calculation requires the annihilation products to fall into thermal equilibrium rapidly with the thermal bath.
- ▶

$$s\text{Hz} \frac{dY_\chi}{dz} = -\langle\sigma v\rangle_{\chi\chi\rightarrow N_{(D)}N_{(D)}} s^2 (Y_\chi^2 - Y_{\chi\text{eq}}^2) \quad (2)$$

# Standard WIMP Calculation?

- ▶ Standard WIMP calculation requires the annihilation products to fall into thermal equilibrium rapidly with the thermal bath.



$$s\text{Hz} \frac{dY_\chi}{dz} = -\langle\sigma v\rangle_{\chi\chi\rightarrow N_{(D)}N_{(D)}} s^2 (Y_\chi^2 - Y_{\chi\text{eq}}^2) \quad (2)$$

- ▶ For naive seesaw model,  $y_\nu \ll 1$ ,  $N$  might deviate from the thermal equilibrium with the thermal bath!

# Standard WIMP Calculation?

- ▶ Standard WIMP calculation requires the annihilation products to fall into thermal equilibrium rapidly with the thermal bath.



$$s\text{Hz} \frac{dY_\chi}{dz} = -\langle\sigma v\rangle_{\chi\chi \rightarrow N_{(D)}N_{(D)}} s^2 (Y_\chi^2 - Y_{\chi\text{eq}}^2) \quad (2)$$

- ▶ For naive seesaw model,  $y_\nu \ll 1$ ,  $N$  might deviate from the thermal equilibrium with the thermal bath!
- ▶ For the pseudo-Dirac sterile neutrinos,  $y_{N_D}$  can  $\sim 0.01$ , however, when  $T < m_{N_D}$ , the effective decay/inverse-decay rate drops rapidly.

# Standard WIMP Calculation?

- ▶ Standard WIMP calculation requires the annihilation products to fall into thermal equilibrium rapidly with the thermal bath.



$$s\text{Hz} \frac{dY_\chi}{dz} = -\langle\sigma v\rangle_{\chi\chi \rightarrow N_{(D)}N_{(D)}} s^2 (Y_\chi^2 - Y_{\chi\text{eq}}^2) \quad (2)$$

- ▶ For naive seesaw model,  $y_\nu \ll 1$ ,  $N$  might deviate from the thermal equilibrium with the thermal bath!
- ▶ For the pseudo-Dirac sterile neutrinos,  $y_{N_D}$  can  $\sim 0.01$ , however, when  $T < m_{N_D}$ , the effective decay/inverse-decay rate drops rapidly.
- ▶ Secluded dark matter.

# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ New progress: the relic abundance of such kind of model.

# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ New progress: the relic abundance of such kind of model.
- ▶ A model independent analysis cannot formulate a complete and reliable  $\langle\sigma v\rangle(T)$  at any temperature, so we rely on a simple model based on Miguel Escudero, et.al, 1607.02373.

# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ New progress: the relic abundance of such kind of model.
- ▶ A model independent analysis cannot formulate a complete and reliable  $\langle\sigma v\rangle(T)$  at any temperature, so we rely on a simple model based on Miguel Escudero, et.al, 1607.02373.
- ▶ A majorana spinor  $\chi$  and a real-scalar  $\phi$  take the minus  $Z_{(2,DM)}$  charge and  $m_\chi < m_\phi$ , so  $\chi$  is the dark matter candidate.

# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ New progress: the relic abundance of such kind of model.
- ▶ A model independent analysis cannot formulate a complete and reliable  $\langle\sigma v\rangle(T)$  at any temperature, so we rely on a simple model based on Miguel Escudero, et.al, 1607.02373.
- ▶ A majorana spinor  $\chi$  and a real-scalar  $\phi$  take the minus  $Z_{(2,DM)}$  charge and  $m_\chi < m_\phi$ , so  $\chi$  is the dark matter candidate.
- ▶  $\chi + \chi \rightarrow N_{(D)} + N_{(D)}$  through the  $\chi N_{(D)}\phi$ -interaction.



# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ New progress: the relic abundance of such kind of model.
- ▶ A model independent analysis cannot formulate a complete and reliable  $\langle\sigma v\rangle(T)$  at any temperature, so we rely on a simple model based on Miguel Escudero, et.al, 1607.02373.
- ▶ A majorana spinor  $\chi$  and a real-scalar  $\phi$  take the minus  $Z_{(2,DM)}$  charge and  $m_\chi < m_\phi$ , so  $\chi$  is the dark matter candidate.
- ▶  $\chi + \chi \rightarrow N_{(D)} + N_{(D)}$  through the  $\chi N_{(D)}\phi$ -interaction.
- ▶  $\phi$  can interact with the Higgs boson through the  $\phi\phi H^\dagger H$  terms.

# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ When the temperature  $T \gg m_{\phi, \chi, N_{(D)}}$ , everything become in thermal equilibrium with the thermal bath through the Higgs $\leftrightarrow\phi$  portal processes.

# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ When the temperature  $T \gg m_{\phi, \chi, N_{(D)}}$ , everything become in thermal equilibrium with the thermal bath through the Higgs $\leftrightarrow\phi$  portal processes.
- ▶ As the temperature  $T$  drops and  $\phi$  decouples,  $N_{(D)}$  and  $\chi$  together decouple from the thermal bath while they are in thermal-equilibrium within themselves.

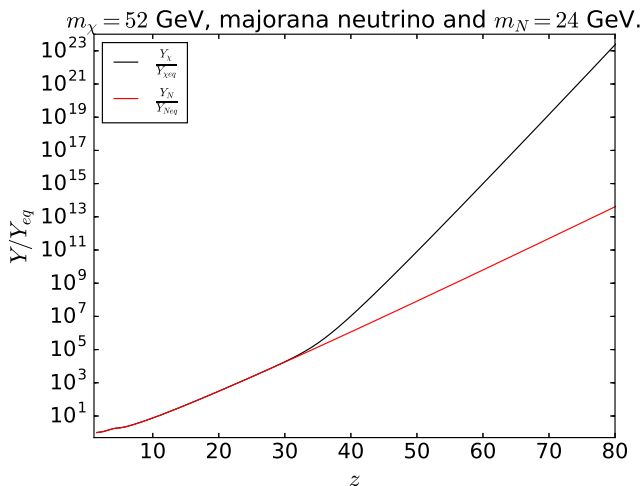
# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ When the temperature  $T \gg m_{\phi, \chi, N_{(D)}}$ , everything become in thermal equilibrium with the thermal bath through the Higgs $\leftrightarrow\phi$  portal processes.
- ▶ As the temperature  $T$  drops and  $\phi$  decouples,  $N_{(D)}$  and  $\chi$  together decouple from the thermal bath while they are in thermal-equilibrium within themselves.
- ▶ Finally,  $N_{(D)}$  and  $\chi$  decouple with each other and  $N_{(D)}$  decays up before the BBN.

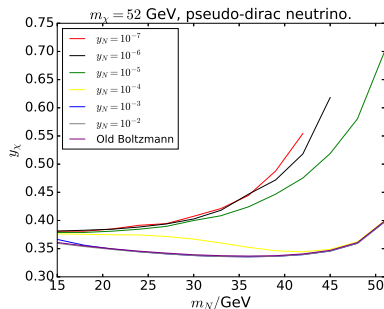
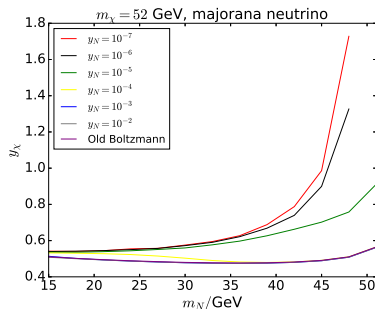
# Introduction to a Simple Sterile-Neutrino Portal Model.

- ▶ When the temperature  $T \gg m_{\phi, \chi, N_{(D)}}$ , everything become in thermal equilibrium with the thermal bath through the Higgs $\leftrightarrow\phi$  portal processes.
- ▶ As the temperature  $T$  drops and  $\phi$  decouples,  $N_{(D)}$  and  $\chi$  together decouple from the thermal bath while they are in thermal-equilibrium within themselves.
- ▶ Finally,  $N_{(D)}$  and  $\chi$  decouple with each other and  $N_{(D)}$  decays up before the BBN.
- ▶ The contribution from the  $W/Z/\gamma^T$  in the thermal bath was estimated according to the method introduced in Phys.Rev.Lett. 117 (2016) no.9, 091801.

- Calculate in a completed Boltzmann equation,



- ▶ In order for a correct relic abundance, the interactions of the  $\chi\chi \rightarrow N_{(D)}N_{(D)}$  should be stronger than the usual standard WIMP calculations!



- ▶ Future plan: to build a complete supersymmetric/nonsupersymmetric model that dark matter  $\rightarrow$  RHN. Explaining neutrino mass spectrum and mixing patterns, leptogenesis, etc....



Thank You!

Thank You!